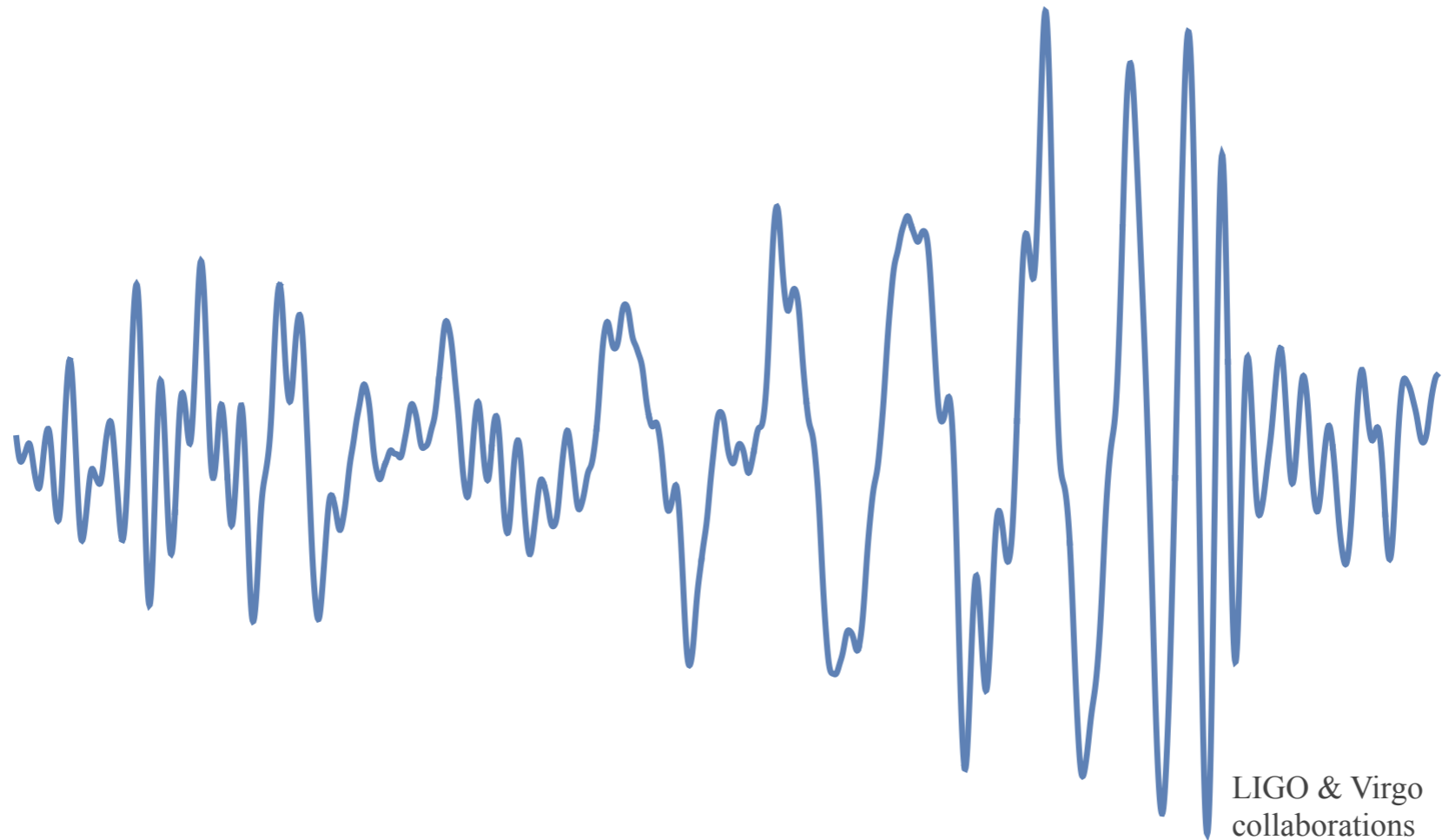


University of Hertfordshire, March 2024

Amplitudes and Waveforms

Donal O'Connell
Edinburgh

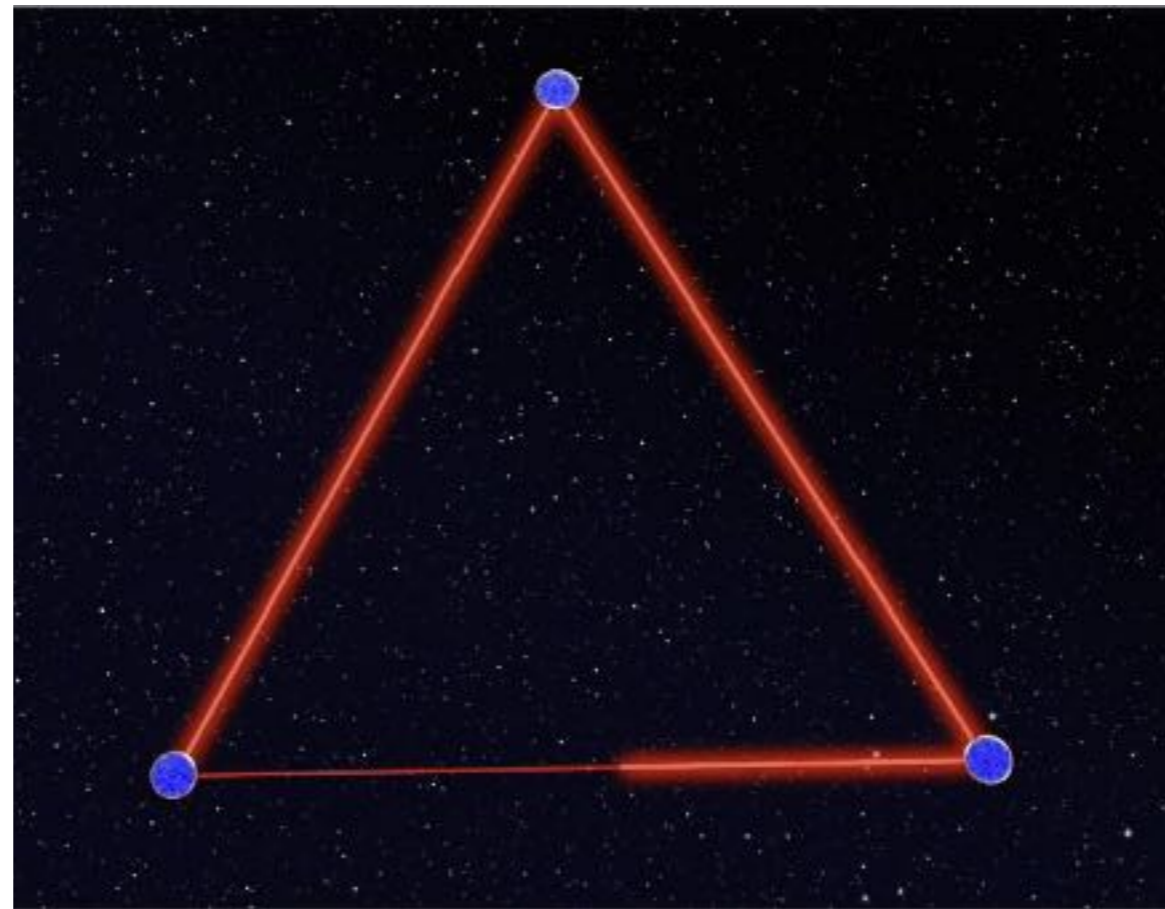
Motivation



Gravity: data rich

Motivation

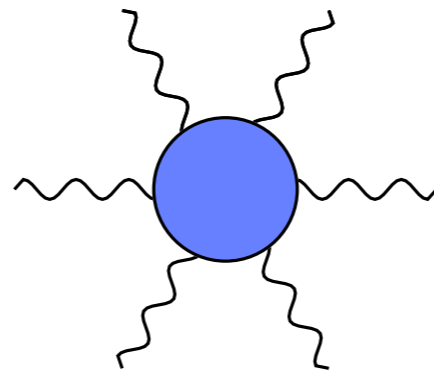
Eg LISA



High-precision gravitational wave theory required

Motivation

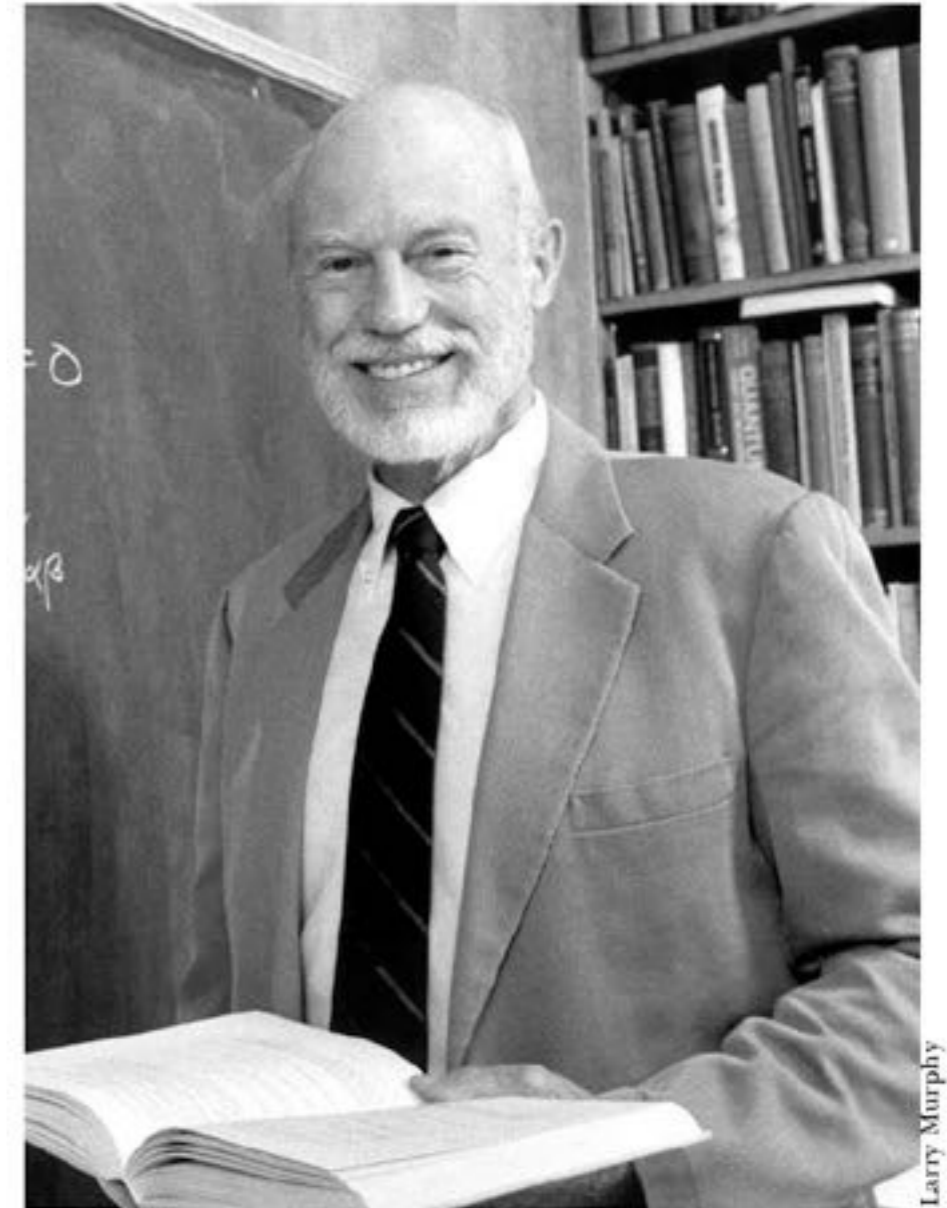
Early inspiral is perturbative



Motivation

Describe method for computing classical observables from amplitudes

“Only observable in (quantum) gravity in asymptotically flat space is the S-matrix”

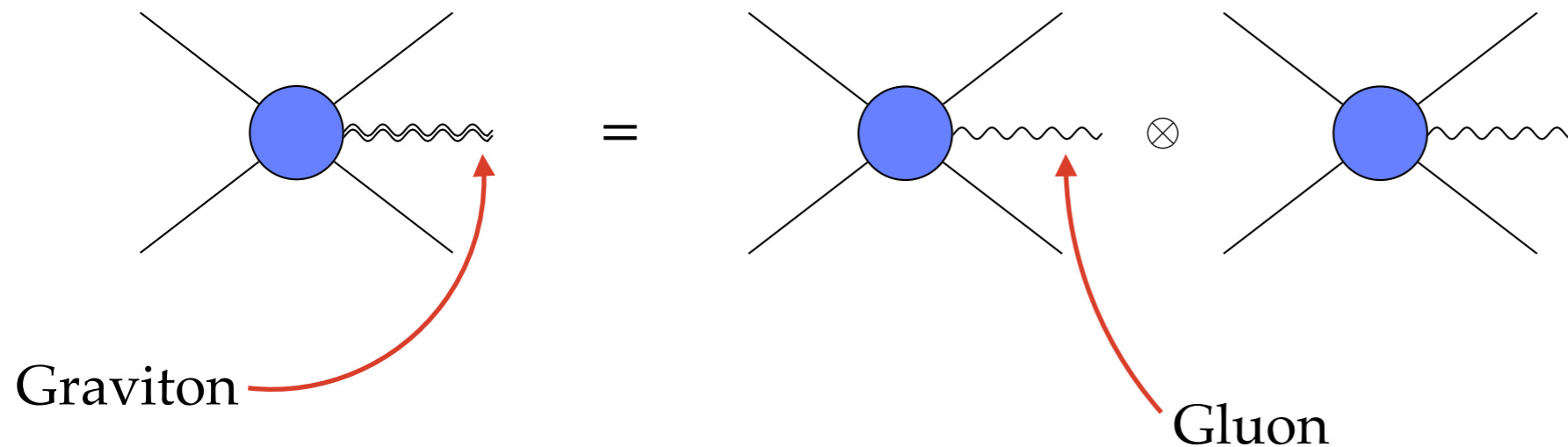


Engel De Wit

Motivation

Amplitudes: different perspective on gravity

The “double copy”:



*Kawai, Lewellen, Tye
Bern, Carrasco, Johansson
Cachazo, He, Yuan*

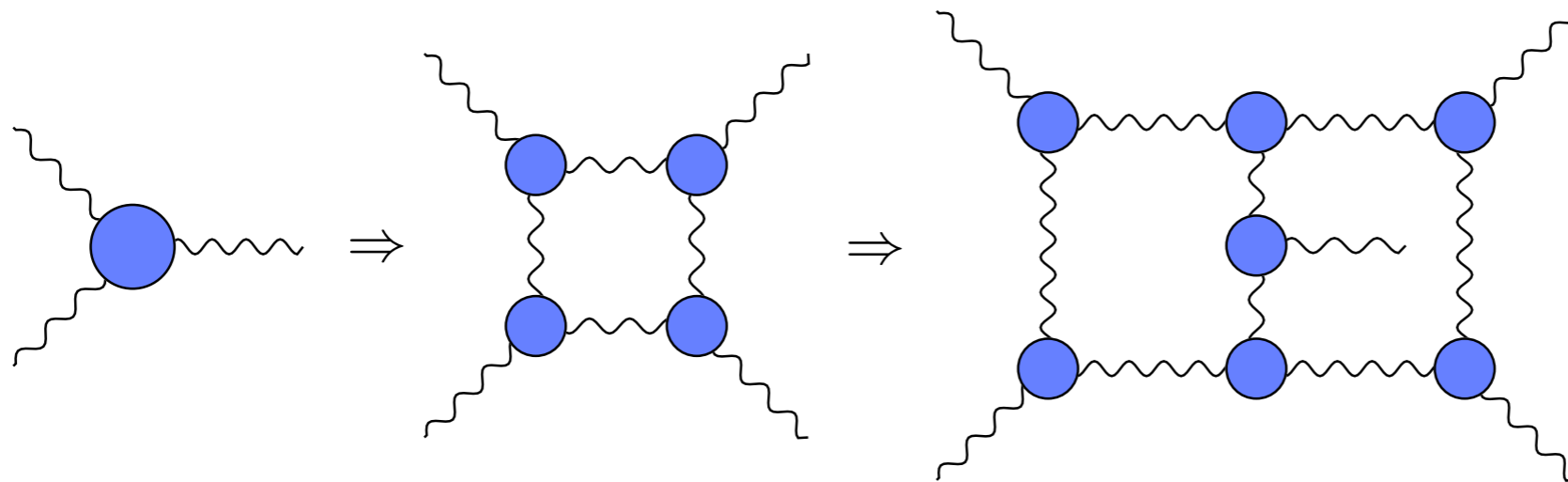
Gravity = Yang-Mills²

Very bizarre from geometric point of view

Motivation

Amplitudes: different perspective on gravity

The unitarity method:



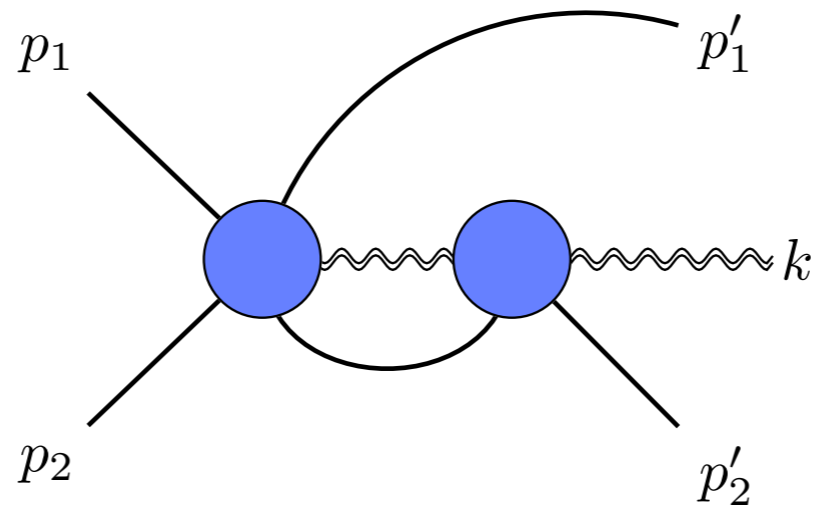
*Bern, Dixon, Dunbar, Kosower
Neill, Rothstein*

Gravitational processes determined by locality / quantum unitarity

At least novel from classical point of view

Motivation

My motivation: *learn about amplitudes*



Waveform \neq cross-section

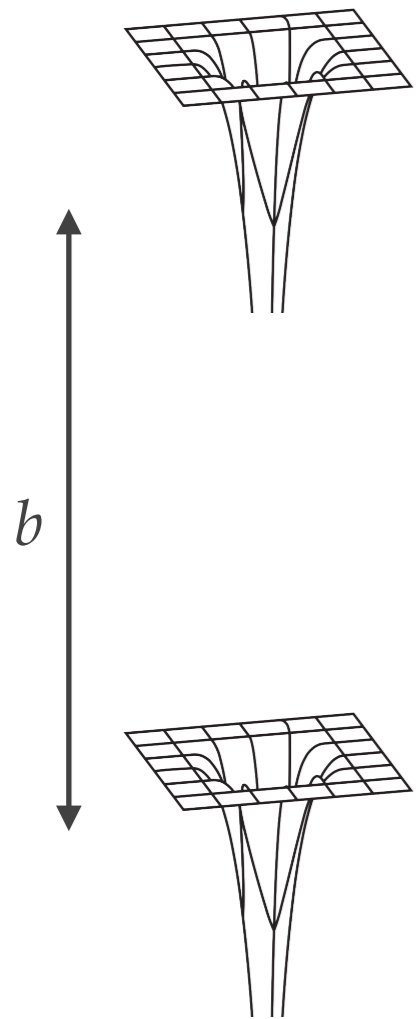
New questions, new challenges

This talk

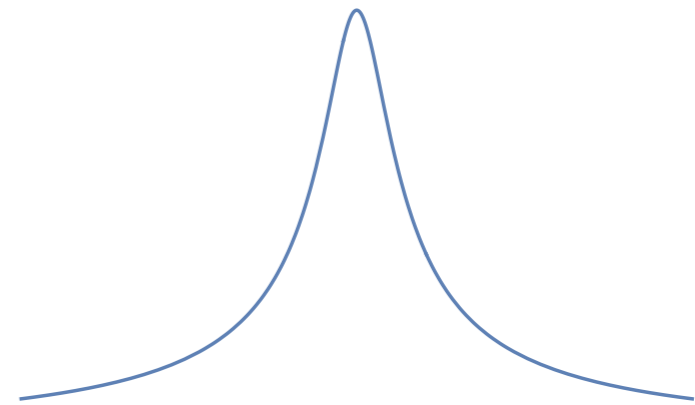
1. Waveforms from amplitudes
2. Waveform @ NLO
 1. In-in expectation
 2. Radiation reaction
 3. Integrals
 4. Surprises
3. Conclusions

Waveforms from Amplitudes

Waves from Amplitudes



$$\text{waveform}(t) =$$



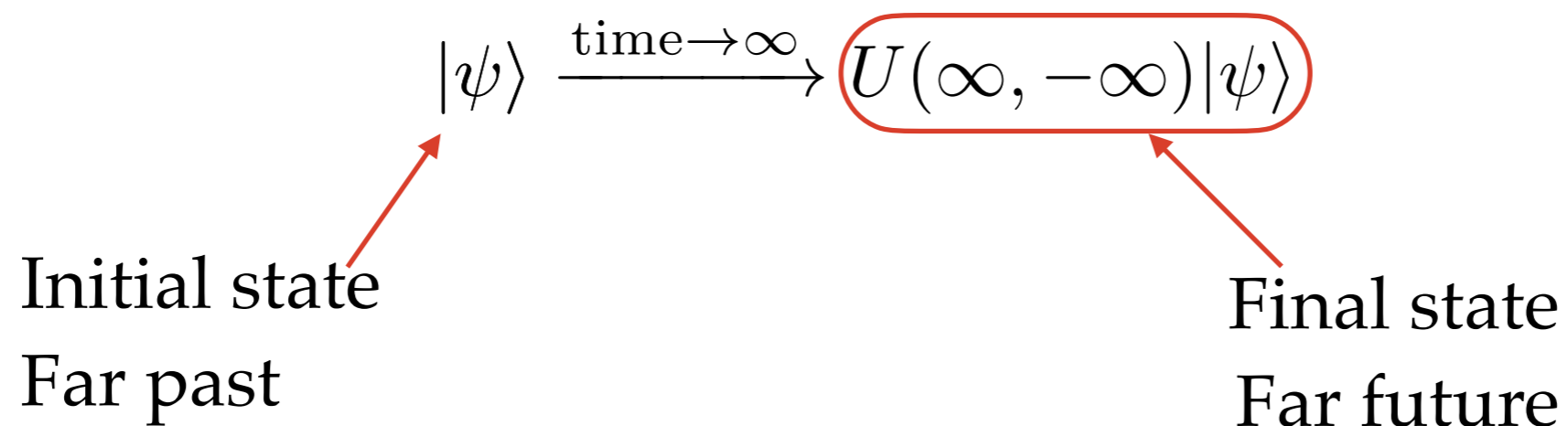
$$= \frac{1}{\text{distance}} \int \text{[diagram of a source point emitting waves]} + \dots$$

Classical point particle approximation: finite size under control

Kosower, Maybee & DOC
Cristofoli, Gonzo, Kosower & DOC

Waveforms & Amplitudes

Amplitudes arise from time evolution



Waveforms & Amplitudes

Amplitudes arise from time evolution

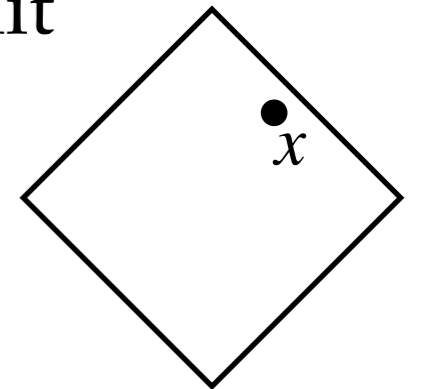
$$|\psi\rangle \xrightarrow{\text{time} \rightarrow \infty} S|\psi\rangle \quad U(-\infty, \infty) = S = 1 + iT$$

Matrix elements of T are the amplitudes

$$\langle q_1 \cdots q_m | T | p_1 \cdots p_n \rangle = \mathcal{A}(p_1 \cdots p_n \rightarrow q_1 \cdots q_m) \delta^4(\text{total momentum})$$

Waveforms & Amplitudes

Measure *expectation* of curvature component in classical limit



Riemann curvature
operator (linear)

$$\text{waveform} \equiv \langle \psi | S^\dagger \mathbb{R} \dots (x) S | \psi \rangle$$

Final state
Amplitudes!

Newman-Penrose scalar Ψ_4
Dominant curvature component
at large distances

$$|\psi\rangle \sim \int \mathcal{N} \mathcal{N} e^{ip_1 \cdot b} |p_1 p_2\rangle$$

Classical initial conditions

E&M analogue useful

$$\text{EM waveform} \equiv \langle \psi | S^\dagger \mathbb{F} \dots (x) S | \psi \rangle$$

Waveforms & Amplitudes

Measure *expectation* of curvature component in classical limit

$$\mathbb{R}....(x) = \partial.\partial.h..(x) \quad \text{Graviton polarisation}$$

$$\text{waveform} = \int \widetilde{d}k [kk \varepsilon\varepsilon e^{-ik \cdot x} \langle \psi | S^\dagger a(k) S | \psi \rangle + \text{c.c.}]$$

$$= i \int \widetilde{d}k kk \varepsilon\varepsilon e^{-ik \cdot x} \langle \psi | a(k) T | \psi \rangle + \dots + \text{c.c.}$$

LO!

Waveforms & Amplitudes

Measure *expectation* of curvature component in classical limit

$$\mathbb{R}....(x) = \partial.\partial.h..(x)$$

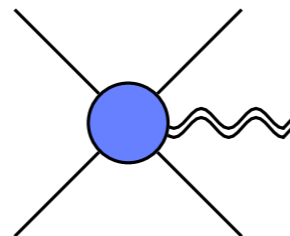
Graviton polarisation

$$\text{waveform} = \int \widetilde{d}k [kk \varepsilon\varepsilon e^{-ik \cdot x} \langle \psi | S^\dagger a(k) S | \psi \rangle + \text{c.c.}]$$

$$= i \int \widetilde{d}k kk \varepsilon\varepsilon e^{-ik \cdot x} \langle \psi | a(k) T | \psi \rangle + \dots + \text{c.c.}$$

LO!

5 point amplitude



Waveforms & Amplitudes

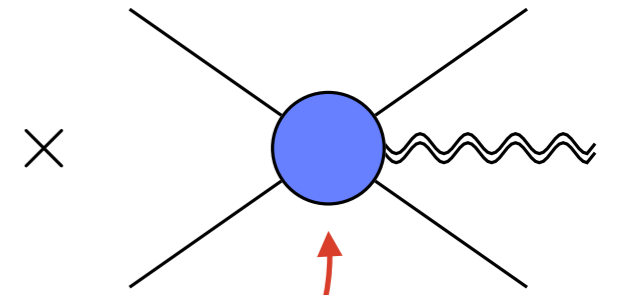
Leading order:

Fourier integral: one bottleneck

$$\Psi_4(\omega) = \frac{1}{\text{distance}} \int d^4 q_1 d^4 q_2 \delta(p_1 \cdot q_1) \delta(p_2 \cdot q_2) \delta^4(k - q_1 - q_2) e^{ib \cdot (q_1 - q_2)}$$

Waveform as a function
of frequency

$$k^\mu = (\omega, \omega \hat{\mathbf{n}})$$



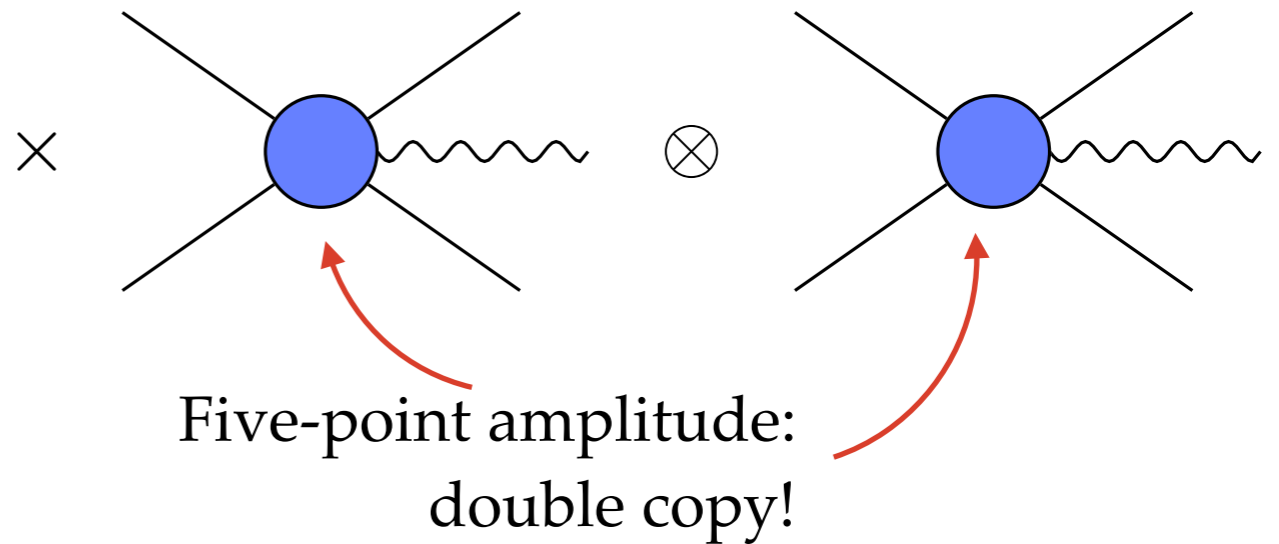
Five-point amplitude:
double copy!

Waveforms & Amplitudes

Leading order:

$$\Psi_4(\omega) = \frac{1}{\text{distance}} \int d^4 q_1 d^4 q_2 \delta(p_1 \cdot q_1) \delta(p_2 \cdot q_2) \delta^4(k - q_1 - q_2) e^{ib \cdot (q_1 - q_2)}$$

LO: Kovacs & Thorne, long ago



Eikonal connection

Seems contrary to intuition:

$$\text{waveform} = \frac{1}{\text{distance}} \int \text{[diagram of a blue circle with four lines and a wavy line]} + \dots$$

One graviton \neq classical field

Classical field: expectation of a coherent state

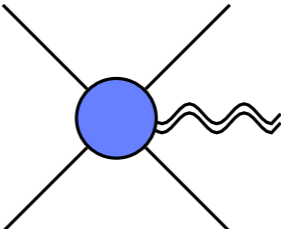
$$\exp \left(\int \tilde{d}k \alpha(k) a^\dagger(k) \right) |0\rangle$$

Waveshape:

Fourier modes of classical field

Eikonal connection

Seems contrary to intuition:

$$\text{waveform} = \frac{1}{\text{distance}} \int \text{diagram} + \dots$$


Problem resolved if amplitude exponentiates in classical region

$$S|\psi\rangle \sim \int \exp\left(i\tilde{\mathcal{M}}_4(x, q)\right) \exp\left(\int \tilde{d}k \left(\mathcal{M}_5(x_1, x_2, k) + \dots\right) a^\dagger(k)\right) |p'_1, p'_2\rangle$$

Generalisation of eikonal exponentiation

Ciafaloni, Colferai, Veneziano

Cristofoli, Gonzo, Moynihan, Ross, Sergola, White, DOC

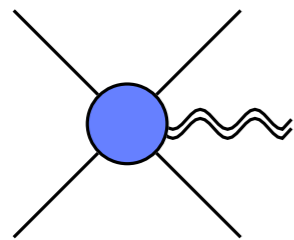
Di Vecchia, Heissenberg, Russo, Veneziano

Bound binaries

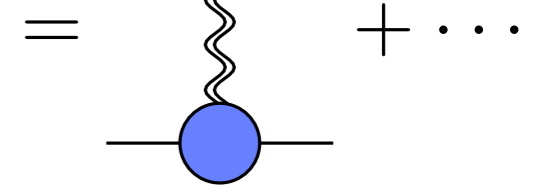
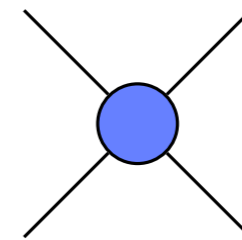
Bound case?

$$\text{Gravitational wave power} = -\frac{d}{dt} \text{Potential}$$

Einstein quadrupole power
+ *corrections*



Newton potential
+ *corrections*



Build EFT, valid in both bound & scattering cases

Match Wilson coefficients to scattering

Neill, Rothstein
Cheung, Rothstein, Solon
Bern, Cheung, Roiban, Shen, Solon, Zeng
Kälin, Porto

Waveform at NLO

Alessio, Bini, Brandhuber, Brown, Bohnenblust, Caron-Huot, Chen, Damour, De Angelis, Di Vecchia, Elkhidir, Giroux, Geralico, Georgoudis, Gowdy, Hannesdottir, Heissenberg, Herderschee, Ita, Krauss, Mizera, Roiban, Russo, Schlenk, Sergola, Teng, Travaglini, Vazquez-Holm, DOC, ..

1. In-in Expectation

$$\text{waveform} = \int \widetilde{d}k [kk \varepsilon\varepsilon e^{-ik \cdot x} \langle \psi | S^\dagger a(k) S | \psi \rangle + \text{c.c.}]$$

Key object to understand

$$\langle \psi | S^\dagger a(k) S | \psi \rangle = \int_{p_i, p'_i} \psi^*(p'_1, p'_2) \psi(p_1, p_2) \langle p'_1, p'_2 | S^\dagger a(k) S | p_1, p_2 \rangle$$

In-in expectation \mathcal{E}

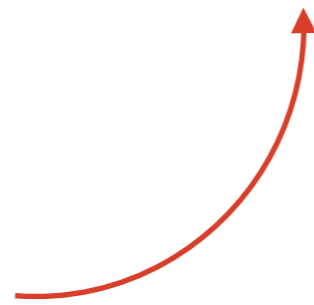
1. In-in Expectation

Expectation closely related to amplitudes

$$\mathcal{E} \delta^4(\text{momentum}) \equiv -i \langle p'_1 p'_2 | S^\dagger a(k) S | p_1 p_2 \rangle$$

$$= \langle p'_1 p'_2 k | T | p_1 p_2 \rangle - i \langle p'_1 p'_2 | T^\dagger a(k) T | p_1 p_2 \rangle$$

LO: tree amplitude



Higher orders: subtraction term



1. In-in Expectation

Define future creation / annihilation operators

$$b(k) \equiv S^\dagger a(k) S$$

$$i\mathcal{A}\delta^4(\text{momentum}) = \langle 0 | b(p'_1) b(p'_2) b(k) a^\dagger(p_1) a^\dagger(p_2) | 0 \rangle$$

$$i\mathcal{E}\delta^4(\text{momentum}) = \langle 0 | a(p'_1) a(p'_2) b(k) a^\dagger(p_1) a^\dagger(p_2) | 0 \rangle$$

Physics: different $i\epsilon$ prescription. In-out vs in-in, Schwinger-Keldysh

Expectation admits classical limit, amplitudes doesn't

1. In-in Expectation

Real parts equal

Treat polarisation vectors as real



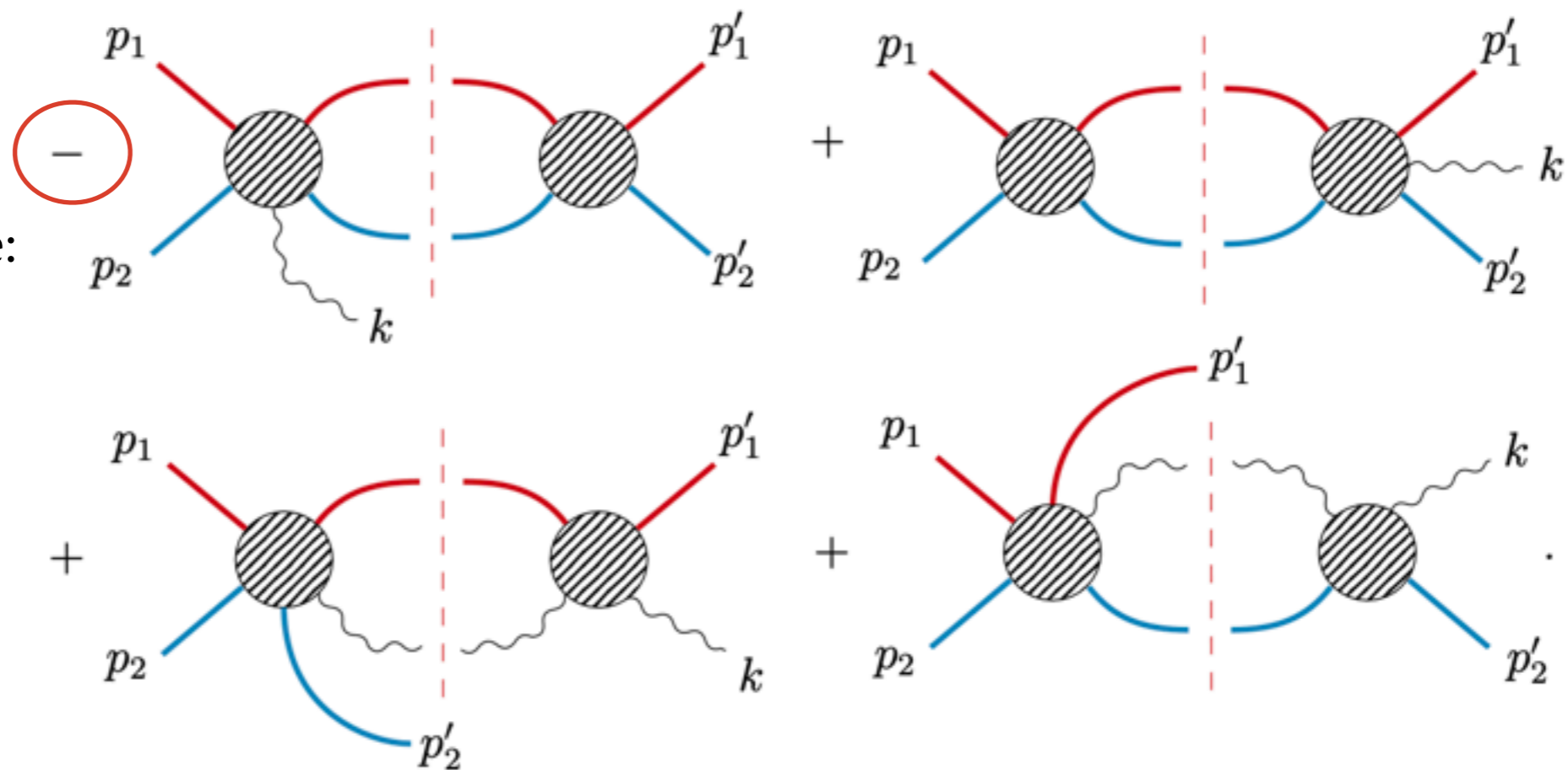
$$\text{Re}' \mathcal{E} = \text{Re}' \mathcal{A}$$

1. In-in Expectation

Im parts differ. At NLO:

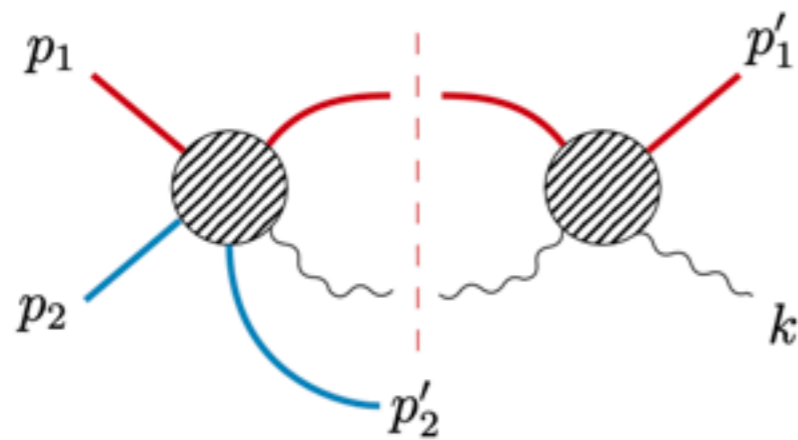
$$2 \text{Im}' \mathcal{E}(p_1 p_2 \rightarrow p'_1 p'_2 k_\eta) \hat{\delta}^D(p_1 + p_2 - p'_1 - p'_2 - k) =$$

Amplitude:
plus sign

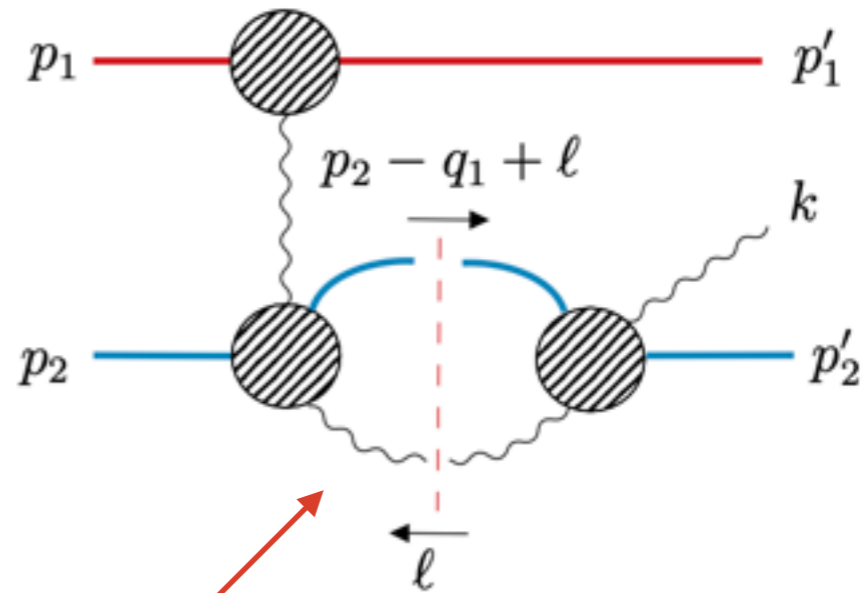


2. Radiation reaction

Focus on E&M



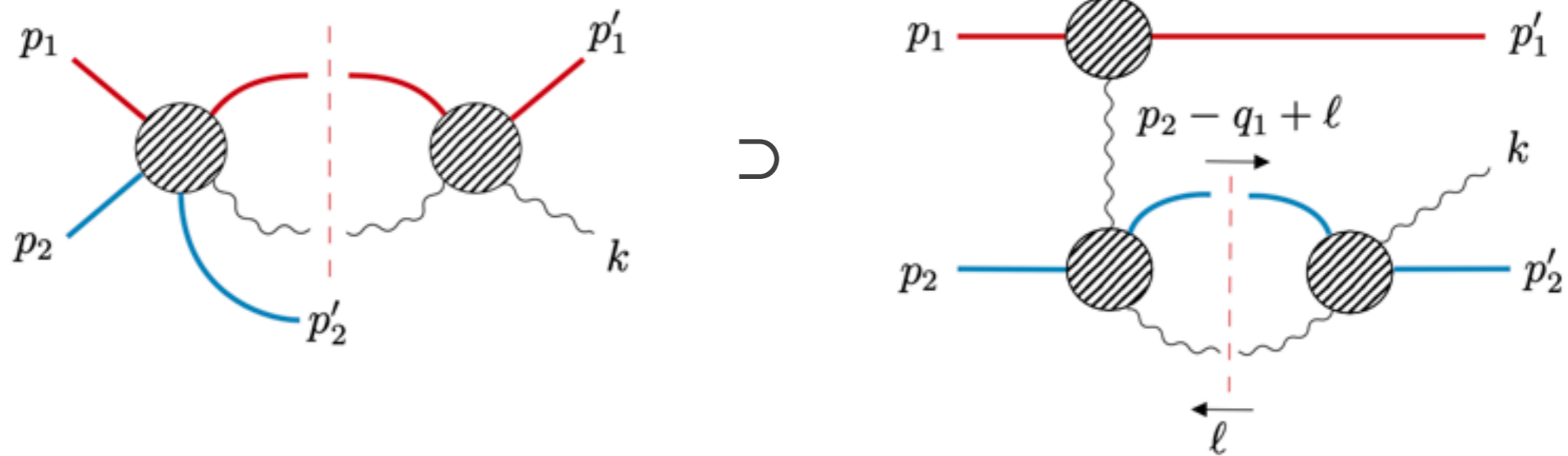
\supset



Particle 2 interacts with its own field

Elkhidir, Sergola, Vazquez-Holm, DOC

2. Radiation reaction



Omit self-field in undergrad EM — divergent, Abraham-Lorentz-Dirac

Amplitudes: textbook one-loop renormalization

Im part well-defined, physical, easily reproduces ALD

Would be interesting to understand in GR, YM

3. Integration

Loop and Fourier integrals

$$\int d^D \ell v^\mu \frac{\partial}{\partial \ell^\mu} \frac{N(\ell)}{(\ell^2 + i\epsilon)(p \cdot \ell + i\epsilon) \cdots} = 0$$

$$\int d^D q v^\mu \frac{\partial}{\partial q^\mu} e^{iq \cdot b} \frac{N(q)}{(q^2 + i\epsilon)(p \cdot q + i\epsilon) \cdots} = 0$$

One loop waveform \sim two loop complexity (in frequency domain)

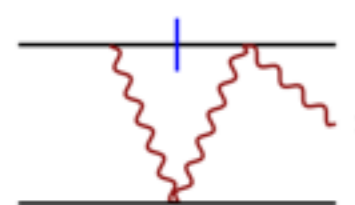
*Matsumoto; Majima, Matsumoto, Takayama; Cacciatori, Mastrolia;
Brunello, Crisanti, Giroux, Mastrolia, Smith; Brunello, De Angelis*

3. Integration

GR waveform currently only numerical

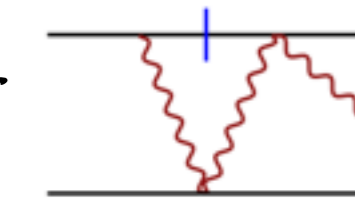
E&M: computed analytically (frequency space) last week!

Brunello & De Angelis



$$= \frac{\pi + 2i \log \frac{w_1 + \sqrt{w_1^2 - q_2^2}}{\sqrt{-q_2^2}}}{8\pi \sqrt{w_1^2 - q_2^2}} + \mathcal{O}(\epsilon^1)$$

Struve function, index -1



$$\mathcal{F} = \frac{i}{16\pi \sqrt{-b^2 p_\infty}} \left\{ z \int_0^\infty dx \left[e^{-z \cosh x} \mathbf{H}_{-1}(z \sqrt{p_\infty} \sinh x) \right] - i \frac{e^{-z \sqrt{1+p_\infty}}}{\sqrt{p_\infty}} \right\}$$

Bessel type

$$z = \frac{w_2 |\mathbf{b}|}{\sqrt{\gamma^2 - 1}}$$

4. Surprises

- ❖ Waveform is IR divergent
 - ❖ But IR divergence is different to divergence of amplitudes

Caron-Huot, Giroux, Hannesdóttir, Mizera

4. Surprises

- ❖ Waveform is IR divergent

- ❖ But IR divergence is different to divergence of amplitudes

Caron-Huot, Giroux, Hannesdóttir, Mizera

- ❖ Post-newtonian (small velocity) expansion compared to classical results (matched multipolar post-Minkowski approach)

- ❖ Agreement found but choice of BMS frame important

- ❖ “Intrinsic” BMS frame: include zero-energy 3-point amplitudes

Venziano, Vilkovisky, Bini, Damour, De Angelis, Geralico, Herderschee, Roiban, Teng

Conclusions

- ❖ Interesting dialogue between amplitudes and classical gravity
- ❖ Waveform constructed from in-in expectation
 - ❖ Close relationship to amplitude
 - ❖ Different pole prescription for some particles
- ❖ Technical challenges being overcome at one loop
- ❖ Conceptual challenges remain: IR divergence, BMS
- ❖ Need to understand bound states